Computational Physics

Lecture-01

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Contents

● Basis Concepts

- Computational Physics aims at **solving physical problems** by means of numerical methods developed in the field of numerical analysis
- In most cases it is not possible to find analytic solutions and one must rely on good approximations.
- In contrast to the much simpler integral
- $\int_{a}^{b}e^{x}\mathrm{d}x=e^{a}-e^{b}% \frac{1}{2\pi}\int_{a}^{b}e^{x}\mathrm{d}x$
- Hence, we have to approximate this integral in such a way that the approximation is accurate enough for our purpose.

• We know from basic mechanics that the time evolution of a frictionless pendulum of mass m and length l in a gravitational field is modeled by the differential equation

$$
\ddot{\theta} + \frac{g}{l}\sin\theta = 0
$$

$$
\theta(t=0) = \theta_0 \text{ and } \dot{\theta}(t=0) = 0
$$

$$
\theta \ll 1 : \sin \theta \approx \theta
$$
 $\ddot{\theta} + \frac{g}{l} \theta = 0$
= $\theta_0 \cos(\omega t)$, $\omega = \sqrt{\frac{g}{l}}$, $T_0 = 2\pi \sqrt{\frac{l}{g}}$

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• We know from basic mechanics that the time evolution of a frictionless pendulum of mass m and length l in a gravitational field is modeled by the differential equation

$$
\ddot{\theta} + \frac{g}{l}\sin\theta = 0
$$

 $\ddot{\theta} = \frac{d\dot{\theta}}{dt} = \frac{d\dot{\theta}}{dt} \frac{d\theta}{dt} = \dot{\theta} \frac{d\dot{\theta}}{dt}$

$$
\theta(t=0) = \theta_0 \text{ and } \dot{\theta}(t=0) = 0
$$

$$
\begin{pmatrix}\n0 \\
0 \\
0 \\
0 \\
0\n\end{pmatrix}
$$

$$
\dot{\theta}\frac{\mathrm{d}\dot{\theta}}{\mathrm{d}\theta} + \frac{g}{l}\sin\theta = 0 \Rightarrow \dot{\theta}\mathrm{d}\dot{\theta} = -\frac{g}{l}\sin\theta\mathrm{d}\theta \Rightarrow \int_0^{\dot{\theta}} \dot{\theta}\mathrm{d}\dot{\theta} = -\frac{g}{l}\int_{\theta_0}^{\theta} \sin\theta\mathrm{d}\theta
$$

• We know from basic mechanics that the time evolution of a frictionless pendulum of mass m and length l in a gravitational field is modeled by the differential equation

$$
\int_0^{\dot{\theta}} \dot{\theta} d\dot{\theta} = -\frac{g}{l} \int_{\theta_0}^{\theta} \sin \theta d\theta
$$

$$
\frac{1}{2}\dot{\theta}^{2} = \frac{g}{l}\left(\cos\theta - \cos\theta_{0}\right) = \frac{g}{l}\left(1 - 2\sin^{2}\frac{\theta}{2} - 1 + 2\sin^{2}\frac{\theta_{0}}{2}\right)
$$

$$
\frac{1}{2}\dot{\theta}^2 = 2\frac{g}{l}\left(\sin^2\frac{\theta_0}{2} - \sin^2\frac{\theta}{2}\right) \Rightarrow \dot{\theta} = 2\sqrt{\frac{g}{l}}\sqrt{\sin^2\frac{\theta_0}{2} - \sin^2\frac{\theta}{2}}
$$

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• We know from basic mechanics that the time evolution of a frictionless pendulum of mass m and length l in a gravitational field is modeled by the differential equation

$$
\dot{\theta} = 2\sqrt{\frac{g}{l}}\sqrt{\sin^2\frac{\theta_0}{2} - \sin^2\frac{\theta_0}{2}} - \sin^2\frac{\theta_0}{2}
$$

$$
\frac{d\theta}{dt} = 2\sqrt{\frac{g}{l}}\sqrt{\sin^2\frac{\theta_0}{2} - \sin^2\frac{\theta_0}{2}} \Rightarrow dt = \frac{1}{2}\sqrt{\frac{l}{g}}\sqrt{\frac{g}{g}}\sqrt{\sin^2\frac{\theta_0}{2} - \sin^2\frac{\theta_0}{2}}
$$

$$
\frac{1}{4}\int_0^T dt = \frac{1}{2}\sqrt{\frac{l}{g}}\int_0^{\theta_0} \frac{d\theta}{\sqrt{\sin^2\frac{\theta_0}{2} - \sin^2\frac{\theta_0}{2}}}
$$

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● A frictionless pendulum

$$
T = 2\sqrt{\frac{l}{g}} \int_0^{\theta_0} \frac{d\theta}{\sqrt{\sin^2 \frac{\theta_0}{2} - \sin^2 \frac{\theta}{2}}}
$$

$$
\begin{array}{c}\n\overline{11}\n\end{array}
$$
\n
$$
\begin{array}{c}\n\overline{11}\n\end{array}
$$

$$
k = \sin\frac{\theta_0}{2} : T = \frac{2}{k} \sqrt{\frac{l}{g}} \int_0^{\theta_0} \frac{d\theta}{\sqrt{1 - \frac{1}{k^2} \sin^2 \frac{\theta}{2}}}
$$

• A frictionless pendulum

$$
T = \frac{2}{k} \sqrt{\frac{l}{g}} \int_0^{\theta_0} \frac{d\theta}{\sqrt{1 - \frac{1}{k^2} \sin^2 \frac{\theta}{2}}}
$$

\n
$$
k \sin \alpha = \sin \frac{\theta}{2}, \qquad 0 \le \sin \frac{\theta}{2} \le \sin \frac{\theta_0}{2} \Rightarrow 0 \le \alpha \le \pi/2
$$

\n
$$
T = 4 \sqrt{\frac{l}{g}} \int_0^{\pi/2} \frac{d\alpha}{\sqrt{1 - k^2 \sin^2 \alpha}} = 4 \sqrt{\frac{l}{g}} F(k, \frac{\pi}{2})
$$

 $F(k, \frac{\pi}{2})$: Complete elliptic integral of the first kind

● A frictionless pendulum

$$
T = 4\sqrt{\frac{l}{g}} \int_0^{\pi/2} \frac{\mathrm{d}\alpha}{\sqrt{1 - k^2 \sin^2 \alpha}} = 4\sqrt{\frac{l}{g}} F(k, \frac{\pi}{2})
$$

Amplitude,

$$
\mathbf{k} = \sin\left(\frac{\theta_0}{2}\right) \qquad \mathbf{F}\left(\mathbf{k}, \frac{\pi}{2}\right) \qquad \text{Period,}
$$

0°	0	1.5708 = $\pi/2$	T ₀	
0°	0	1.5708 = $\pi/2$	T ₀	
10°	0.0872	1.5738	1.0019 T ₀	
45°	0.3827	1.6336	1.0400 T ₀	
90°	0.7071	1.8541	1.1804 T ₀	
135°	0.9234	2.4003	1.5281 T ₀	
178°	0.99985	5.4349	3.5236 T ₀	
180°	1	∞	∞	∞

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● A frictionless pendulum

$$
T = 4\sqrt{\frac{l}{g}} \int_0^{\pi/2} \frac{d\alpha}{\sqrt{1 - k^2 \sin^2 \alpha}} = 4\sqrt{\frac{l}{g}} F(k, \frac{\pi}{2})
$$

$$
k, \frac{\pi}{2}) = \frac{\pi}{2} \sum_0^{\infty} \left[\frac{(2n)!}{2^{2n} (n!)^2} \right]^2 k^{2n} = \frac{\pi}{2} \sum_0^N \left[\frac{(2n)!}{2^{2n} (n!)^2} \right]^2 k^{2n} + R_N(k)
$$

$$
R_N(k): \text{ Truncation Error}
$$

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• A frictionless pendulum

$$
T = 2\pi \sqrt{\frac{l}{g}} \sum_{0}^{\infty} \left[\frac{(2n)!}{2^{2n}(n!)^2} \right]^2 k^{2n} = T_0 \sum_{0}^{\infty} \left[\frac{(2n)!}{2^{2n}(n!)^2} \right]^2 k^{2n}
$$

$$
\int_{0}^{\infty} \left[\frac{(2n)!}{2^{2n}(n!)^2} \right]^2 k^{2n} = \frac{T}{T_0} \Rightarrow k = ? \text{ or } \theta_0 = ?
$$

However, such an approach would be very inefficient due to two reasons:

- We are confronted with the impossibility of finding analytically the roots of a polynomial of order N>4.
- At which value should we truncate the power series?

- Classify the occurring errors based on the structure every numerical routine.
	- * Input-Error □ Roundoff errors in input data
	- * Algorithmic-Error

 \Box Roundoff errors during evaluation □ Methodological errors due to mathematical approximations

* Output-Error □ Roundoff errors

• Roundoff errors

Since every number is stored in a computer using a finite number of digits, we have to truncate every non-terminating number at some point.

$$
\frac{2}{3} = 0.666666666666666\dots
$$
\n
$$
\xrightarrow{\text{forced}} 0.66666666667 \text{ truncated}}
$$
\n
$$
\text{floating-point form } FL\left(\frac{2}{3}\right) = 0.66666666667
$$

This has the consequence: $FL(\sqrt{3})FL(\sqrt{3}) \neq FL(\sqrt{3}\sqrt{3}) = 3$

● Roundoff errors

$$
\epsilon_a = | \ y - \bar{y}
$$

True value
$$
y
$$

Approximate value \bar{y}

$*$ Relative error

$$
\epsilon_r = \left| \frac{y - \bar{y}}{y} \right| = \frac{\epsilon_a}{|y|}
$$

In most applications, the relative error is more significant.

● Roundoff errors

True value
$$
y = 0.d_1d_2d_3\cdots d_kd_{k+1}\cdots \times 10^n
$$

Approximate value $\bar{y} = 0.d_1d_2d_3\cdots d_k \times 10^n$ $d_1 \neq 0$

$$
\epsilon_r = \left| \frac{0.d_1d_2d_3\cdots d_kd_{k+1}\cdots \times 10^n - 0.d_1d_2d_3\cdots d_k \times 10^n}{0.d_1d_2d_3\cdots d_kd_{k+1}\cdots \times 10^n} \right|
$$

$$
\epsilon_r = \left| \frac{0.d_{k+1}d_{k+2}\cdots \times 10^{-k}}{0.d_1d_2d_3\cdots d_kd_{k+1}\cdots} \right| = \left| \frac{0.d_{k+1}d_{k+2}\cdots}{0.d_1d_2d_3\cdots} \right| 10^{-k} < \frac{1}{0.1}10^{-k} = 10^{-k+1}
$$

 $0.d_{k+1}d_{k+2}\cdots < 1$ $d_1 \neq 0: 0.d_1d_2d_3\cdots \geq 0.1,$

• Roundoff errors

* Particular care is required when two nearly identical numbers are subtracted or when a large number is divided by a small number. increase in the relative error.

 $*$ In such cases the roundoff error will increase dramatically.

* We note that it might be necessary to avoid such operations in our aim to design an algorithm which is required to produce reasonable results.

• Roundoff errors

* The machine-number is smallest positive number which can be added to another number, such that a change in the result is observed.

$$
\eta = \min_{\delta} \{ \delta > 0 | 1 + \delta > 1 \}
$$

 $*$ A typical value for **double-precision** in FORTRAN or C is

$$
\eta \approx 10^{-16}
$$

• Methodological Errors

* A methodological error is introduced whenever a complex mathematical expression is replaced by an approximate, or

we have to deal with expressions we cannot evaluate analytically.

$$
\sum_{0}^{\infty} \left[\frac{(2n)!}{2^{2n}(n!)^2} \right]^2 k^{2n} \approx \frac{\pi}{2} \sum_{0}^{N} \left[\frac{(2n)!}{2^{2n}(n!)^2} \right]^2 k^{2n}
$$

$$
f'(x_0) = \frac{df}{dx}\Big|_{x_0} = \lim_{h \to 0} \frac{f(x_0 + h) - f(x_0)}{h} \xrightarrow{\text{Finite difference}} \frac{df}{dx}\Big|_{x_0} \approx \frac{f(x_0 + h) - f(x_0)}{h}
$$