Computational Physics

Lecture-01

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Contents

• Basis Concepts

- Computational Physics aims at **solving physical problems** by means of **numerical methods developed** in the field of numerical analysis
- In most cases it **is not possible to find analytic solutions** and $\int_{a}^{b} e^{-x^{2}} dx$ one must rely on good approximations.
- In contrast to the much simpler integral
- $\int_{a}^{b} e^{x} \mathrm{d}x = e^{a} e^{b}$
- Hence, we have to approximate this integral in such a way that $\int_{a}^{b} e^{-x^{2}} dx \approx \cdots$ the approximation is accurate enough for our purpose.

• We know from basic mechanics that the time evolution of a frictionless pendulum of mass m and length l in a gravitational field is modeled by the differential equation

$$\ddot{\theta} + \frac{g}{l}\sin\theta = 0$$

$$\theta(t=0) = \theta_0$$
 and $\dot{\theta}(t=0) = 0$



$$\theta \ll 1 : \sin \theta \approx \theta \qquad \qquad \ddot{\theta} + \frac{g}{l}\theta = 0$$
$$\theta = \theta_0 \cos(\omega t), \qquad \omega = \sqrt{\frac{g}{l}}, \qquad T_0 = 2\pi \sqrt{\frac{l}{g}}$$

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• We know from basic mechanics that the time evolution of a frictionless pendulum of mass m and length l in a gravitational field is modeled by the differential equation

$$\ddot{\theta} + \frac{g}{l}\sin\theta = 0$$

 $\ddot{\theta} = \frac{\mathrm{d}\dot{\theta}}{\mathrm{d}t} = \frac{\mathrm{d}\dot{\theta}}{\mathrm{d}\theta}\frac{\mathrm{d}\theta}{\mathrm{d}t} = \dot{\theta}\frac{\mathrm{d}\dot{\theta}}{\mathrm{d}\theta}$

$$\theta(t=0) = \theta_0$$
 and $\dot{\theta}(t=0) = 0$

$$\dot{\theta}\frac{\mathrm{d}\dot{\theta}}{\mathrm{d}\theta} + \frac{g}{l}\sin\theta = 0 \Rightarrow \dot{\theta}\mathrm{d}\dot{\theta} = -\frac{g}{l}\sin\theta\mathrm{d}\theta \Rightarrow \int_{0}^{\dot{\theta}}\dot{\theta}\mathrm{d}\dot{\theta} = -\frac{g}{l}\int_{\theta_{0}}^{\theta}\sin\theta\mathrm{d}\theta$$

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• We know from basic mechanics that the time evolution of a frictionless pendulum of mass m and length l in a gravitational field is modeled by the differential equation

$$\int_{0}^{\dot{\theta}} \dot{\theta} \mathrm{d}\dot{\theta} = -\frac{g}{l} \int_{\theta_{0}}^{\theta} \sin\theta \mathrm{d}\theta$$

$$\frac{1}{2}\dot{\theta}^2 = \frac{g}{l}\left(\cos\theta - \cos\theta_0\right) = \frac{g}{l}\left(1 - 2\sin^2\frac{\theta}{2} - 1 + 2\sin^2\frac{\theta_0}{2}\right)$$

$$\frac{1}{2}\dot{\theta}^2 = 2\frac{g}{l}\left(\sin^2\frac{\theta_0}{2} - \sin^2\frac{\theta}{2}\right) \Rightarrow \dot{\theta} = 2\sqrt{\frac{g}{l}}\sqrt{\sin^2\frac{\theta_0}{2} - \sin^2\frac{\theta}{2}}$$

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• We know from basic mechanics that the time evolution of a frictionless pendulum of mass m and length l in a gravitational field is modeled by the differential equation

$$\dot{\theta} = 2\sqrt{\frac{g}{l}}\sqrt{\sin^2\frac{\theta_0}{2}} - \sin^2\frac{\theta}{2}$$
$$\frac{\mathrm{d}\theta}{\mathrm{d}t} = 2\sqrt{\frac{g}{l}}\sqrt{\sin^2\frac{\theta_0}{2}} - \sin^2\frac{\theta}{2} \Rightarrow \mathrm{d}t = \frac{1}{2}\sqrt{\frac{l}{g}}\frac{\mathrm{d}\theta}{\sqrt{\sin^2\frac{\theta_0}{2}} - \sin^2\frac{\theta}{2}}$$
$$\frac{1}{4}\int_0^T \mathrm{d}t = \frac{1}{2}\sqrt{\frac{l}{g}}\int_0^{\theta_0}\frac{\mathrm{d}\theta}{\sqrt{\sin^2\frac{\theta_0}{2}} - \sin^2\frac{\theta}{2}}$$

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• A frictionless pendulum

$$T = 2\sqrt{\frac{l}{g}} \int_0^{\theta_0} \frac{\mathrm{d}\theta}{\sqrt{\sin^2 \frac{\theta_0}{2} - \sin^2 \frac{\theta}{2}}}$$



$$k = \sin \frac{\theta_0}{2}: \quad T = \frac{2}{k} \sqrt{\frac{l}{g}} \int_0^{\theta_0} \frac{\mathrm{d}\theta}{\sqrt{1 - \frac{1}{k^2} \sin^2 \frac{\theta}{2}}}$$

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• A frictionless pendulum

$$T = \frac{2}{k} \sqrt{\frac{l}{g}} \int_{0}^{\theta_{0}} \frac{\mathrm{d}\theta}{\sqrt{1 - \frac{1}{k^{2}} \sin^{2} \frac{\theta}{2}}}$$
$$k \sin \alpha = \sin \frac{\theta}{2}, \qquad 0 \le \sin \frac{\theta}{2} \le \sin \frac{\theta_{0}}{2} \Rightarrow 0 \le \alpha \le \pi/2$$
$$T = 4 \sqrt{\frac{l}{g}} \int_{0}^{\pi/2} \frac{\mathrm{d}\alpha}{\sqrt{1 - k^{2} \sin^{2} \alpha}} = 4 \sqrt{\frac{l}{g}} F(k, \frac{\pi}{2})$$

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 $F(k, \frac{\pi}{2})$: Complete elliptic integral of the first kind

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• A frictionless pendulum

$$T = 4\sqrt{\frac{l}{g}} \int_{0}^{\pi/2} \frac{d\alpha}{\sqrt{1 - k^2 \sin^2 \alpha}} = 4\sqrt{\frac{l}{g}} F(k, \frac{\pi}{2})$$

$$\begin{array}{c|c} \textbf{Amplitude,} & \textbf{k} = \sin\left(\frac{\theta_0}{2}\right) & F\left(\textbf{k}, \frac{\pi}{2}\right) & \textbf{Period,} \\ \hline \theta_0 & 0 & 1.5708 = \pi/2 & T_0 \\ \hline 0^\circ & 0 & 1.5708 = \pi/2 & T_0 \\ \hline 10^\circ & 0.0872 & 1.5738 & 1.0019 T \\ 45^\circ & 0.3827 & 1.6336 & 1.0400 T \\ 90^\circ & 0.7071 & 1.8541 & 1.1804 T \\ 135^\circ & 0.9234 & 2.4003 & 1.5281 T \\ 178^\circ & 0.99985 & 5.4349 & 3.5236 T \\ 179^\circ & 0.99996 & 5.2660 & 4.6002 T \\ 180^\circ & 1 & \infty & \infty \end{array}$$

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• A frictionless pendulum

$$T = 4\sqrt{\frac{l}{g}} \int_{0}^{\pi/2} \frac{\mathrm{d}\alpha}{\sqrt{1 - k^{2} \sin^{2} \alpha}} = 4\sqrt{\frac{l}{g}} F(k, \frac{\pi}{2})$$

$$\int_{g}^{\pi/2} \frac{\mathrm{d}\alpha}{\sqrt{1 - k^{2} \sin^{2} \alpha}} = 4\sqrt{\frac{l}{g}} F(k, \frac{\pi}{2})$$

$$F(k, \frac{\pi}{2}) = \frac{\pi}{2} \sum_{0}^{\infty} \left[\frac{(2n)!}{2^{2n}(n!)^{2}}\right]^{2} k^{2n} = \frac{\pi}{2} \sum_{0}^{N} \left[\frac{(2n)!}{2^{2n}(n!)^{2}}\right]^{2} k^{2n} + R_{N}(k)$$

$$R_{N}(k): \text{ Truncation Error}$$

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• A frictionless pendulum

However, such an approach would be very inefficient due to two reasons:

- We are confronted with the **impossibility of finding analytically the roots** of a polynomial of order *N>4*.
- At which value should we **truncate the power series**?

- Classify the occurring errors based on the structure every numerical routine.
 - * **Input-Error**□ Roundoff errors in input data
 - * Algorithmic-Error
 - Roundoff errors during evaluation
 Methodological errors due to mathematical approximations
 - * **Output-Error** □ Roundoff errors



• Roundoff errors

Since every number is stored in a **computer using a finite number of digits**, we have to **truncate** every non-terminating number at some point.

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This has the consequence: $FL(\sqrt{3})FL(\sqrt{3}) \neq FL(\sqrt{3}\sqrt{3}) = 3$

• Roundoff errors

*	Abso	lute	error
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$$\epsilon_a = \mid y - \bar{y}$$

True value y
Approximate value \bar{y}

* Relative error

$$\epsilon_r = \left| \frac{y - \bar{y}}{y} \right| = \frac{\epsilon_a}{|y|}$$

In most applications, the relative error is more significant.

	y	\bar{y}	ϵ_a	ϵ_r
(1)	0.1	0.09	0.01	0.1
(2)	1000.0	999.99	0.01	0.00001

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• Roundoff errors

True value
$$y = 0.d_1d_2d_3\cdots d_kd_{k+1}\cdots \times 10^n$$

Approximate value $\bar{y} = 0.d_1d_2d_3\cdots d_k \times 10^n$ $d_1 \neq 0$

$$\epsilon_r = \left| \frac{0.d_1 d_2 d_3 \cdots d_k d_{k+1} \cdots \times 10^n - 0.d_1 d_2 d_3 \cdots d_k \times 10^n}{0.d_1 d_2 d_3 \cdots d_k d_{k+1} \cdots \times 10^n} \right|$$

$$\epsilon_r = \left| \frac{0.d_{k+1}d_{k+2}\cdots \times 10^{-k}}{0.d_1d_2d_3\cdots d_kd_{k+1}\cdots} \right| = \left| \frac{0.d_{k+1}d_{k+2}\cdots}{0.d_1d_2d_3\cdots} \right| 10^{-k} < \frac{1}{0.1}10^{-k} = 10^{-k+1}$$

 $d_1 \neq 0: 0.d_1 d_2 d_3 \cdots \geq 0.1, \qquad 0.d_{k+1} d_{k+2} \cdots < 1$

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• Roundoff errors

* Particular care is required when two nearly identical numbers are subtracted or when a large number is divided by a small number. increase in the relative error.

* In such cases the **roundoff error will increase** dramatically.

* We note that it might be necessary to avoid such operations in our aim to design an algorithm which is required to produce reasonable results.

• Roundoff errors

* The **machine-number** is smallest positive number which can be added to another number, such that a change in the result is observed.

$$\eta = \min_{\delta} \{\delta > 0 | 1 + \delta > 1\}$$

* A typical value for **double-precision** in FORTRAN or C is

$$\eta \approx 10^{-16}$$

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• Methodological Errors

* A **methodological error** is introduced whenever a **complex mathematical expression is replaced by an approximate**, or

we have to deal with expressions we cannot evaluate analytically.

$$\sum_{0}^{\infty} \left[\frac{(2n)!}{2^{2n}(n!)^2} \right]^2 k^{2n} \approx \frac{\pi}{2} \sum_{0}^{N} \left[\frac{(2n)!}{2^{2n}(n!)^2} \right]^2 k^{2n}$$
$$f'(x_0) = \left. \frac{\mathrm{d}f}{\mathrm{d}x} \right|_{x_0} = \lim_{h \to 0} \frac{f(x_0 + h) - f(x_0)}{h} \xrightarrow{\text{Finite difference}} \frac{\mathrm{d}f}{\mathrm{d}x} \right|_{x_0} \approx \frac{f(x_0 + h) - f(x_0)}{h}$$

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